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Electron-neutral scattering cross sections for CO₂

Advanced Program in Plasma Science and Engineering - APPLAuSE

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Introduction

Conversion of carbon dioxide (CO₂) using plasma technology is considered a great scientific challenge. A detailed understanding of the first step in the process – the CO₂ dissociation by plasma – is essential for its control and optimization [1]. To address this problem, a comprehensive knowledge of the electron kinetics is required. Accordingly, a reliable quantification of electron-impact cross sections and a critical evaluation of available dissociation cross sections are needed. This work offers a set of electron-neutral scattering cross sections from ground-state CO₂, to be published on the IST-LISBON database with LXCat, and evaluates and recommends cross sections for providing the rate coefficients for electron-impact CO₂ dissociation.

Results and discussion

The proposed swarm-derived complete and consistent set of electron-neutral scattering cross sections from ground-state CO₂ includes 17 cross sections.

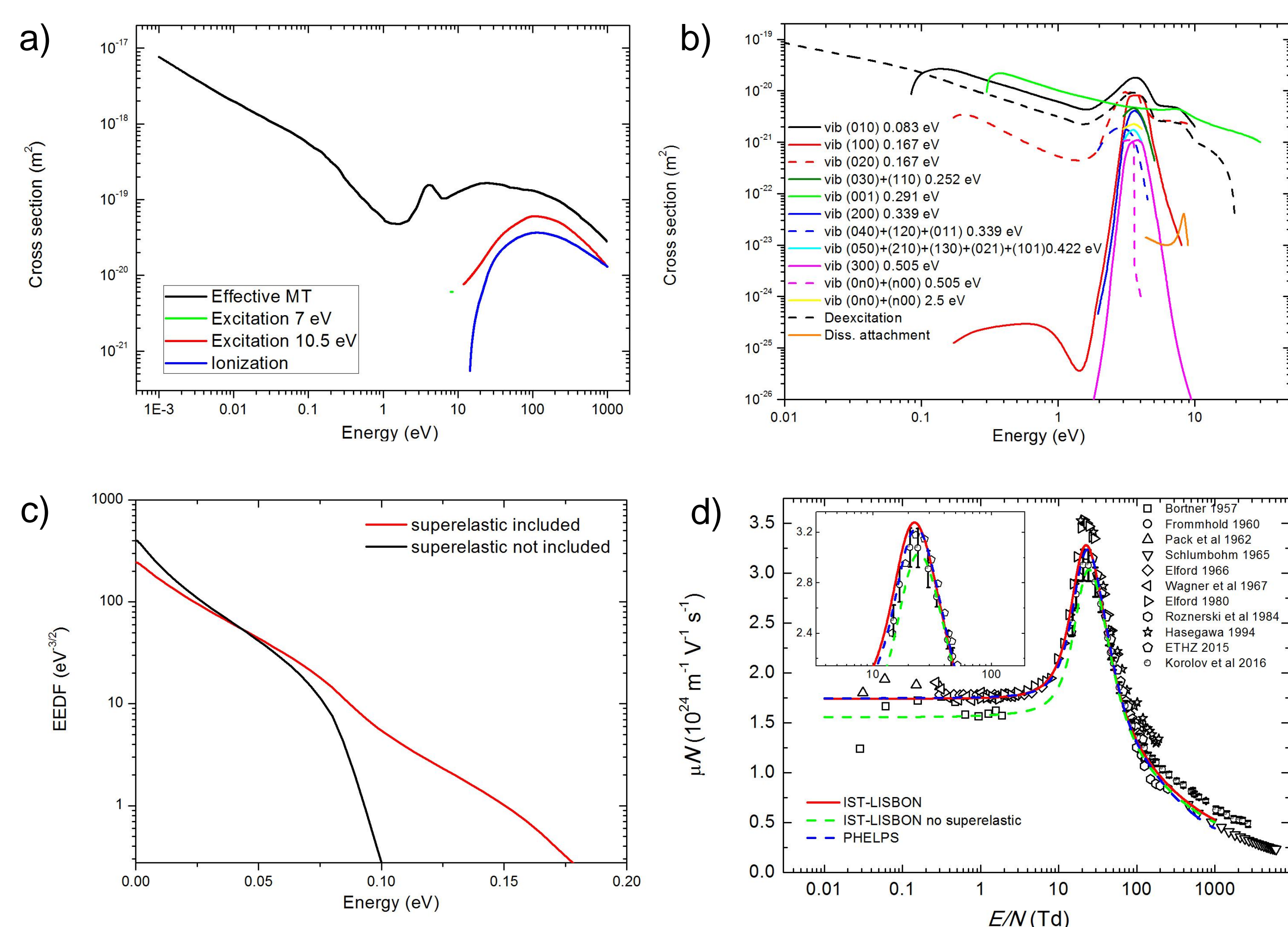


Fig.1. Summary of the proposed CO₂ cross section set, as a function of the electron kinetic energy a) and b); EEDF obtained with the proposed set for 1 Td c); measured and calculated reduced electron mobility as a function of the reduced electric field d).

The set was compiled mostly from Phelps [2] and includes the cross sections defined up to 1000 eV, describing dissociative attachment, effective momentum transfer, eight vibrational excitation energy losses (corresponding either to the excitation of individual levels or groups of vibrational levels), superelastic collisions with the CO₂(010) vibrational state, excitation of two groups of electronic states and and ionization (see table 1).

Table 1. Summary of the processes considered in the cross section set proposed.

	Heavy-species products	Configuration of final CO ₂ state	Threshold [eV]
Effective momentum-transfer	CO ₂ (v ₀)	(000)	
Dissociative attachment	CO+O ⁻		
Vibrational excitation	CO ₂ (v ₁)	(010)	0.083
Superelastic deexcitation	CO ₂ (v ₀)	(000)	
Vibrational excitation	CO ₂ (v _{2a})	(020)	0.167
Vibrational excitation	CO ₂ (v _{2b})	(100)	0.167
Vibrational excitation	CO ₂ (v ₃)	(030)+(110)	0.252
Vibrational excitation	CO ₂ (v ₄)	(001)	0.291
Vibrational excitation	CO ₂ (v _{5a})	(200)	0.339
Vibrational excitation	CO ₂ (v _{5b})	(040)+(120)+(011)	0.339
Vibrational excitation	CO ₂ (v ₆)	(050)+(210)+(130)+(021)+(101)	0.442
Vibrational excitation	CO ₂ (v _{7a})	(300)	0.505
Vibrational excitation	CO ₂ (v _{7b})	(0n0)+(n00)	0.505
Vibrational excitation	CO ₂ (v ₈)	(0n0)+(n00)	2.500
Electronic excitation	CO ₂ (e ₁)		7.0
Electronic excitation	CO ₂ (e ₁)		10.5
Total ionization	CO ₂ ⁺		13.3

The cross sections are summarized in Figs. 1a) and b). A small number modifications was made in regard to the original set by Phelps. First, superelastic collisions with the first level of the bending mode were included and should be considered as an integral part of the set. As a matter of fact, due to its low energy threshold (0.08 eV), this level presents a non-negligible population even in thermal conditions and influences the electron kinetics, specially at low values of the reduced electric field. In the original set [2] the electronic excitation with the threshold at 10.5 eV and the ionization cross sections are limited to 100 eV kinetic energy. The cross section corresponding to the aforementioned electronic excitation was extended up to 1000 eV, while the ionization cross section was replaced by the total ionization cross section from [3]. Considering the extension of the inelastic cross section for $u > 100$ eV, the effective momentum transfer cross section is modified in the same energy region. Finally, the effective momentum transfer cross section was slightly increased for electron energies below 1eV, in order to somehow compensate for the additional gain of energy associated with the superelastic collisions.

When this complete set of cross sections is used as input data in a two-term Boltzmann code it yields calculated swarm parameters in agreement with available swarm measurements, as discussed in [4,5]. The present calculations show that the energy gained in the superelastic process is relevant for low ($E/N > 10$ Td) reduced electric fields (Fig. 1c, d).

The cross section set presented in this work will be soon available in the IST LISBON database with LXCat (www.lxcat.net). At present, the IST-LISBON database includes complete and consistent sets of electron scattering cross sections for argon, helium, nitrogen, oxygen, hydrogen and methane, concluded by the Group of Gas Discharges and Gaseous Electronics with Instituto de Plasmas e Fusão Nuclear, Instituto Superior Técnico, Lisboa, Portugal.

References

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